

Transient Behavior of Grounding Systems Embedded in Different Earth Structures

Fani E. Asimakopoulou, Vassiliki T. Kontargyri, Ioannis F. Gonos and Ioannis A. Stathopoulos
School of Electrical and Computer Engineering, National Technical University of Athens,
Athens, GR 15780, Greece

Abstract— A grounding system has to ensure personal safety with regard to the developed step and touch voltages, avert damage to properties and equipment, and achieve reliability of the power system. In this paper, the software package PC OPERA has been used for the simulation of the transient behavior of a driven rod embedded into different earth structures. This software package is able to calculate electromagnetic fields by using the Finite Element Method in order to estimate the corresponding vector and scalar potentials. Moreover, the influence of the voltage waveform on the behavior of the rod has been studied. The results demonstrate the contribution of the earth structure to the developed potential around the rod and the flexibility of the software package when simulating transient phenomena.

Index Terms-- earth structure, FEM, grounding systems, potential distribution, transient behavior

I. INTRODUCTION

THE grounding system comprises an essential part of the protection of any structure (power system, building, etc.), dissipates the stroke current into the soil and reduces the damages to personnel as well as to electrical and electronic equipment. The commonly used grounding structures are single horizontal grounding wires, vertical rods, ring conductors or a suitable combination of the above mentioned structures. The behavior of a grounding system, when subjected to impulse currents, differs from its steady state behavior. Many works have been published, describing methodologies for analyzing the transient behavior of grounding systems. Earlier models proposed by Bewley [1], Sunde [2], Gupta [3], which were based on analytical methods, suffer from simplification assumptions. However, numerical methods overcome geometrical limitations and can be applied on modeling complex grounding systems. Meliopoulos [4], [5], Geri [6], Otero [7] have suggested a circuit approach for studying the transient behavior of grounding systems. Verma [8], Mazzetti [9], Velazquez [10], Grcev [11], [12], Liu [13], [14], have adopted and evolved the transmission line approach. However, the most accurate method is the electromagnetic field approach which solves full Maxwell's equations. Finite Element Method [15]-[17], Method of Moments [18], Boundary Element Method [19], [20] or hybrid methods [21] can be used for solving electromagnetic field problems in grounding systems.

II. SIMULATION PROGRAM

Opera-2d is a software package for electromagnetic field

analysis, which solves a wide range of electromagnetic and electrostatic applications in both XY and axisymmetric coordinate systems. This package uses the Finite Element Method to obtain solutions to partial differential equations (Poisson's, Helmholtz and Diffusion equations) that cannot be solved by analytic methods. Opera-2d can also solve electric field problems in models with conductivity and permittivity. Additionally, non-linear materials can be modeled by this program.

Since much information is required before the analysis has been performed, data entry is carried out using a powerful interactive pre-processor. Using the graphical interaction within the pre-processor, the space is divided into a contiguous set of triangular elements.

Once the model has been prepared, the solution is achieved using a suitable analysis module. Several modules exist for analysis of the different types of electromagnetic excitation conditions, e.g. static, transient, steady state. The analysis program iteratively determines the correct solution, including non-linear parameters, if these are modelled.

The result may then be examined using a versatile interactive postprocessor. As with the pre-processor, this is predominantly controlled by interaction through a graphical menu system. Many system variables are available for examination, including potentials, currents, fields, forces, temperature. Numerical errors due to non-successful mesh definition are also analysed, so that the mesh can be refined to achieve the required accuracy [22], [23].

The transient analysis program (Opera-2d/TR) solves eddy current models where the driving currents or voltages are changing in time in a predetermined way. It can analyze the response to multiple drive functions, including a dc background field, skin effect, nonlinear materials, in both XY and axisymmetric coordinate systems [22], [23].

Opera-2d/TR solves the vector diffusion equation with the magnetic vector potential as the unknown variable. It is formed from equations

$$\nabla \times \mathbf{H} = \mathbf{J} \quad (1)$$

$$\mathbf{B} = \mu(\mathbf{H} - \mathbf{H}_c) \quad (2)$$

$$\mathbf{J} = \sigma \mathbf{E} \quad (3)$$

$$\mathbf{B} = \nabla \times \mathbf{A} \quad (4)$$

and the integration of

$$\nabla \times \mathbf{E} = \frac{\partial \mathbf{B}}{\partial t} \quad (5)$$

where \mathbf{H} is the magnetic field strength, \mathbf{J} is the current density, \mathbf{B} is the magnetic flux density, \mathbf{A} is the magnetic vector potential, \mathbf{E} is the electric field strength, μ is the permeability and σ is the conductivity.

Combining the (1)-(5), the following equation has been deduced:

$$\nabla \times \left(\frac{1}{\mu} \nabla \times \mathbf{A} - \mathbf{H}_c \right) = \mathbf{J}_s - \sigma \frac{\partial \mathbf{A}}{\partial t} \quad (6)$$

in which the current density has been split into the source current, \mathbf{J}_s and the induced current, $\sigma \frac{\partial \mathbf{A}}{\partial t}$. In two dimensions, only the z components of \mathbf{A} and \mathbf{J}_s exist. Equation (6) can be simplified to:

$$-\nabla \cdot \frac{1}{\mu} \nabla \times \mathbf{A}_z - (\nabla \times \mathbf{H}_c)_z = \mathbf{J}_s - \sigma \frac{\partial \mathbf{A}_z}{\partial t} \quad (7)$$

In addition, Opera-2d/LD option solves lossy dielectric models [22], [23]. It models the behavior of devices consisting of materials with both conductive and dielectric properties under steady-state and transient conditions. In transient models, the finite element formulation is based on the assumption that inductive effects are negligible in semiconducting dielectric models, leading to the quasi-static electric formulation, the governing equations of which are [25]:

$$\nabla \cdot \left(\sigma \nabla V(P, t) + \frac{\partial}{\partial t} (\epsilon \nabla V(P, t)) \right) = 0 \quad (8)$$

$$\nabla \cdot (\epsilon \nabla V(P, 0)) = 0 \quad (9)$$

where P denotes the space coordinates, t denotes time, V is the electric scalar potential and ϵ is the permittivity.

III. DRIVEN ROD IN A MULTI-LAYER EARTH

The grounding systems consisting of a 3m length and 17.2mm diameter single driven rod buried in different soil types is simulated by using Opera-2d. The soil model adopted in this paper is similar to the one suggested by [26], [27] and is presented in Fig.1. According to this model, the soil consists of two concentric hemispheres of radii $r_1=5\text{m}$, $r_2=10\text{m}$. In Table I the values of each zone's soil resistivity and permittivity for the examined cases (1, 2, 3 and 4) are shown. These values are used for the simulations.

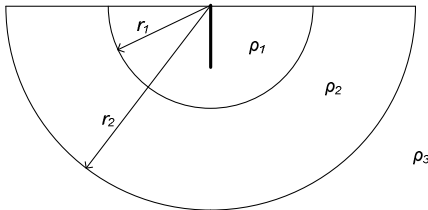


Fig. 1. The geometry of the problem.

TABLE I THE SOIL PARAMETERS USED FOR THE SIMULATIONS				
	Case 1	Case 2	Case 3	Case 4
ρ_1 [Ωm]	50	2000	1000	50
ρ_2 [Ωm]	400	400	50	1000
ρ_3 [Ωm]	1000	100	1000	200
ϵ_1	41	3.5	9	41
ϵ_2	16	16	41	9
ϵ_3	9	36	9	20
R [Ω]	37.70	623.28	324.30	34.52

IV. SIMULATION RESULTS

The arrangement is axisymmetric. Mesh generation is of essential importance for simulation by using Opera-2d. Various meshes have been examined in order to define the most appropriate for the problem. A more dense mesh has been used on the common surface of different materials and in the area around the grounding rod. Fig 2 displays the mesh within the examined regions.

Moreover, the determination of the boundary conditions is equally important. Therefore, the boundary conditions have been carefully selected. In our Opera-2d simulation, the current flow equation was first solved providing the boundary conditions for the electrostatic problem, while the soil resistivity and the relative permittivity (ϵ_r) were defined as presented in Table I.

The double exponential function described by (10) is considered as stimulus of the grounding system:

$$V = V_0 (e^{-at} - e^{-bt}) \quad (10)$$

The values of the parameters V_0 , a , b are presented in Table II.

TABLE II PARAMETERS OF DOUBLE-EXPONENTIAL VOLTAGE		
	1 st Waveform	2 nd Waveform
V_0 [kV]	1.031	1.104
a [μsec^{-1}]	27000	7924
b [μsec^{-1}]	5600000	400109

For the simulations two voltage waveforms are being used and are presented in Fig. 3. It should be noted that the computed values are normalized to 1kV peak value of the induced voltage and proportionally larger for higher voltages.

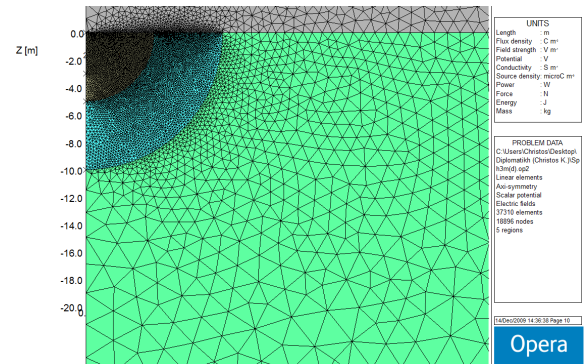


Fig. 2. The Opera-2d model with the mesh for Case 1.

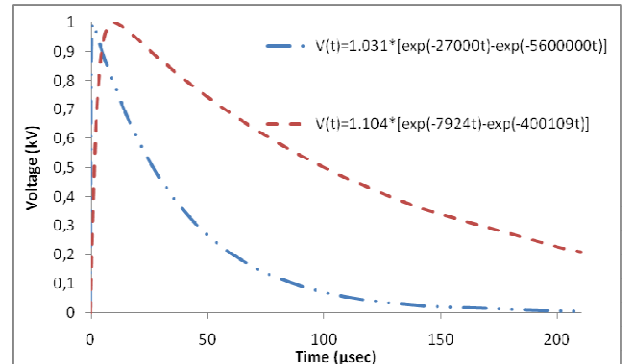


Fig. 3. Voltage stimuli of the grounding system.

In Figs. 4-7 the distribution of the potential on the soil surface is presented for each case at $t=1\mu\text{s}$ when the rod is subjected to the 1st waveform.

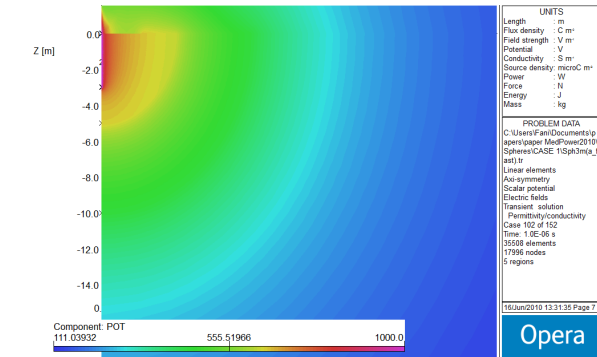


Fig. 4. Filled zone contours of potential for Case 1 up to 20m from the electrode at $t=1\mu\text{s}$.

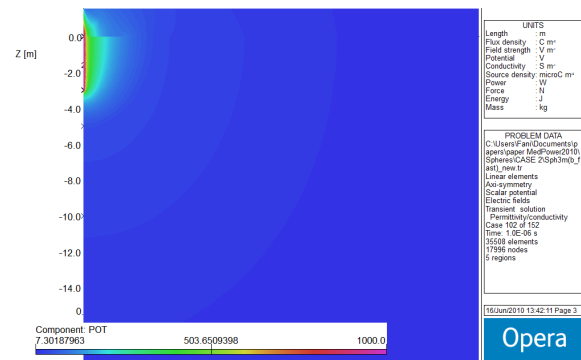


Fig. 5. Filled zone contours of potential for Case 2 up to 20m from the electrode at $t=1\mu\text{s}$.

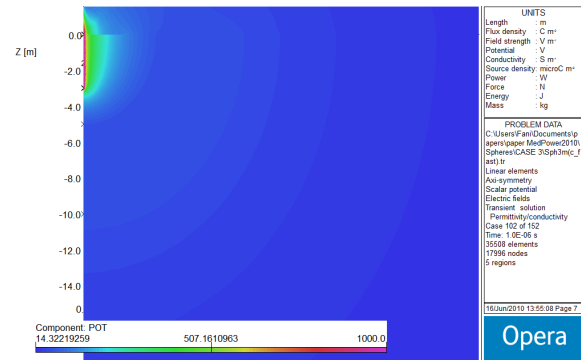


Fig. 6. Filled zone contours of potential for Case 3 up to 20m from the electrode at $t=1\mu\text{s}$.

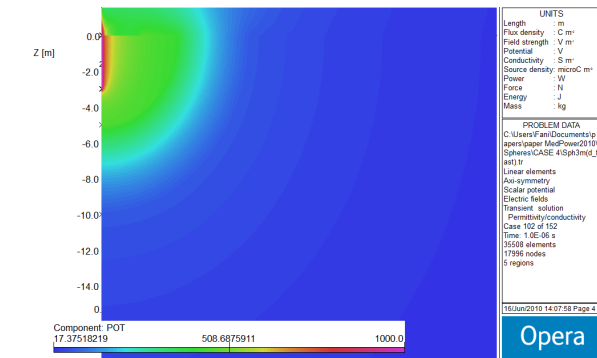


Fig. 7. Filled zone contours of potential for Case 4 up to 20m from the electrode at $t=1\mu\text{s}$.

V. ESTIMATION OF THE SURFACE POTENTIAL

In Figs. 8-15 the time varying distribution of the potential on the soil up to a distance of 50m from the electrode for the examined cases for fast and slow impulse waveforms is presented.

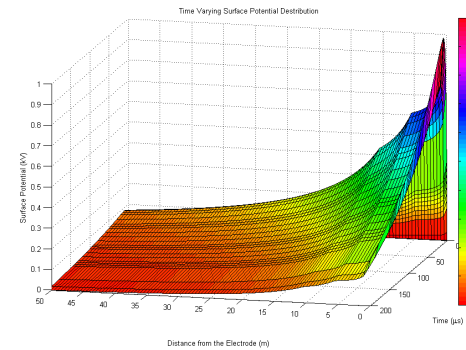


Fig. 8. The time varying surface potential versus time and the horizontal distance from the rod for Case 1 and the 1st Voltage Waveform.

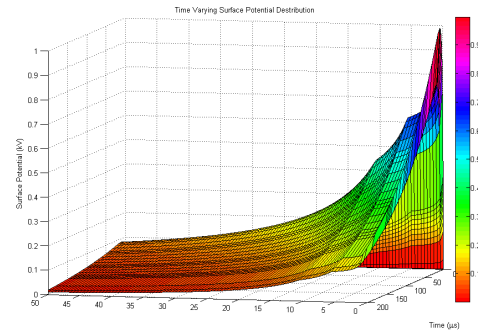


Fig. 9. The time varying surface potential versus time and the horizontal distance from the rod for Case 1 and the 2st Voltage Waveform.

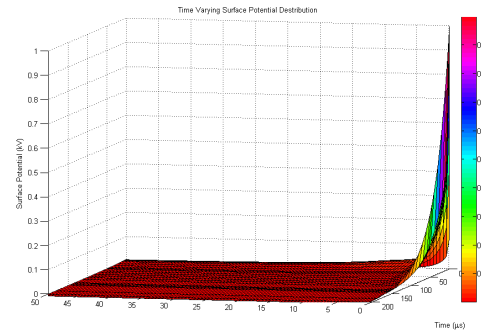


Fig. 10. The time varying surface potential versus time and the horizontal distance from the rod for Case 2 and the 1st Voltage Waveform.

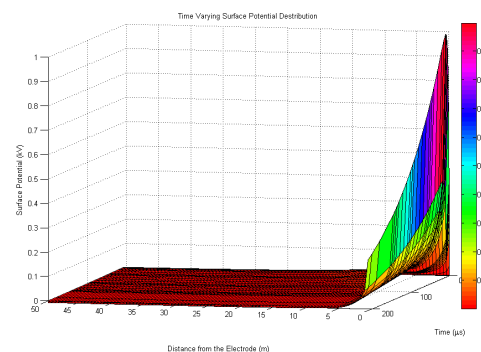


Fig. 11. The time varying surface potential versus time and the horizontal distance from the rod for Case 2 and the 2nd Voltage Waveform.

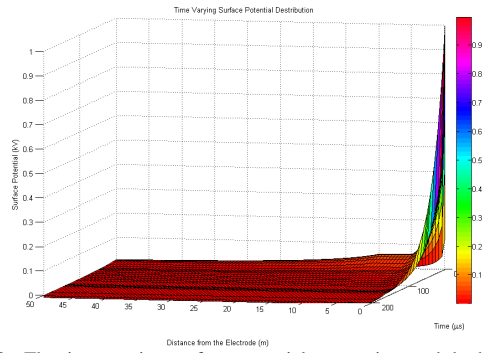


Fig. 12. The time varying surface potential versus time and the horizontal distance from the rod for Case 3 and the 1st Voltage Waveform.

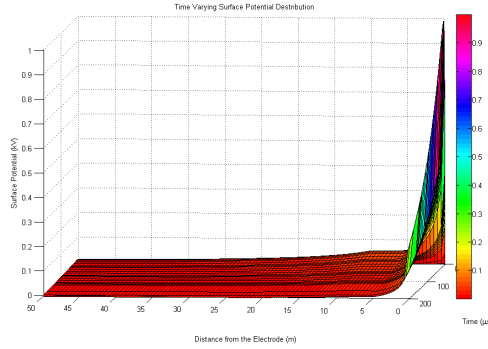


Fig. 13. The time varying surface potential versus time and the horizontal distance from the rod for Case 3 and the 2nd Voltage Waveform.

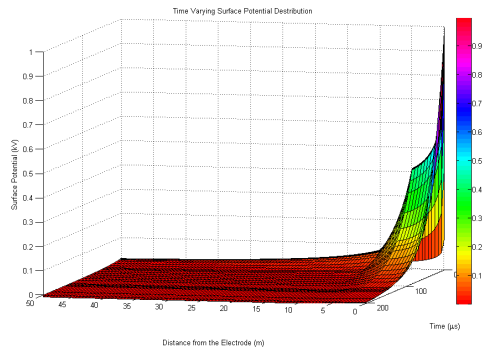


Fig. 14. The time varying surface potential versus time and the horizontal distance from the rod for Case 4 and the 1st Voltage Waveform.

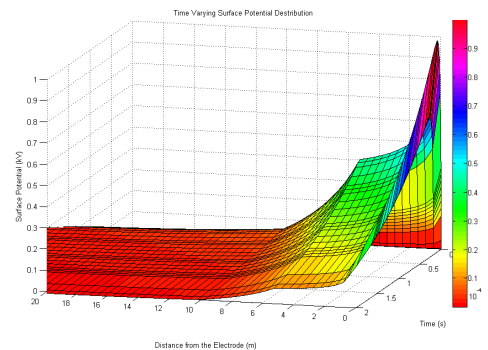


Fig. 15. The time varying surface potential versus time and the horizontal distance from the rod for Case 4 and the 2nd Voltage Waveform.

In Fig. 16 the potential distribution on the soil surface is presented for all the examined cases at the moment of peak voltage of the fast waveform.

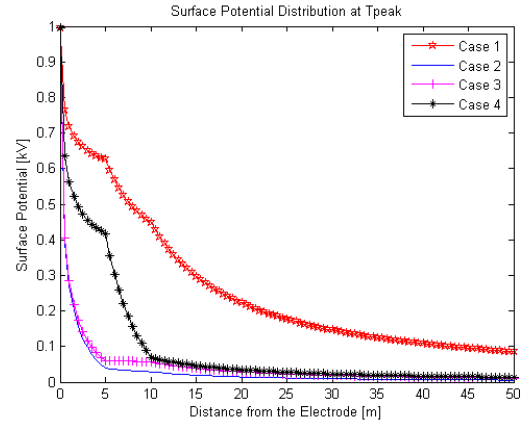


Fig. 16. Surface potential distribution at T_{peak} for all cases.

For all the examined cases, it has been observed that at a distance of 5m and 10m, where the soil properties change, the rate of the potential's drop also changes. In Cases 2 and 3, where the region closer to the electrode is characterized by a lower value of the soil resistivity in comparison to the second region, the potential drop is steeper than in Cases 1 and 4. Therefore, in that zone dangerous step and touch voltages are developed.

Comparing Case 1 to Case 4, the influence of the third (outer) zone on the potential distribution can be examined. In Case 1 where the resistivity increases as the distance from the electrode increases, the potential drop is softer. On the other hand, in Case 4 the potential drops faster, especially as we are moving from the inner zone to the second one. This fact, also, leads to higher step and touch voltages.

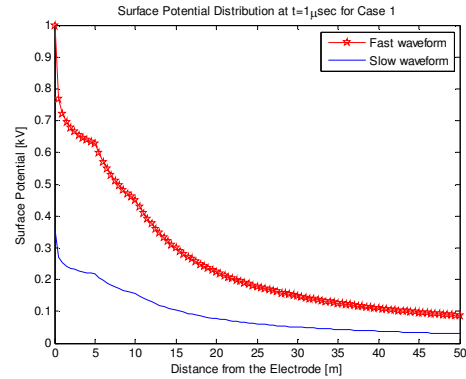


Fig. 17. Surface potential distribution up to a distance of 50m for Case 1 at $t=1 \mu\text{sec}$ under fast and slow impulse waveforms.

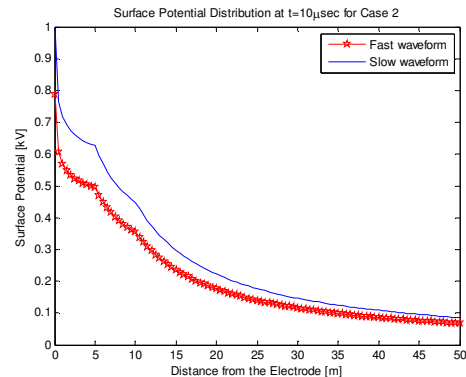


Fig. 18. Surface potential distribution up to a distance of 50m for Case 1 at $t=10 \mu\text{sec}$ under fast and slow impulse waveforms.

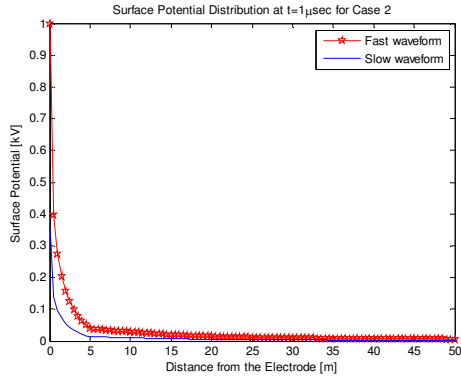


Fig. 19. Surface potential distribution up to a distance of 50m for Case 2 at $t=1\ \mu\text{sec}$ under fast and slow impulse waveforms.

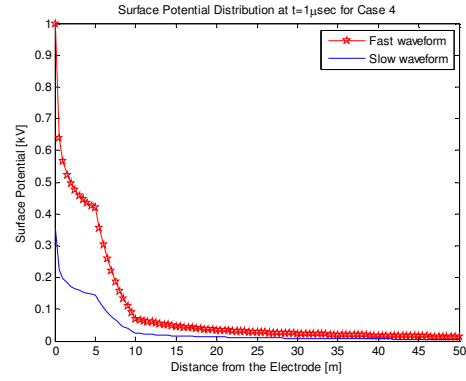


Fig. 23. Surface potential distribution up to a distance of 50m for Case 4 at $t=1\ \mu\text{sec}$ under fast and slow impulse waveforms.

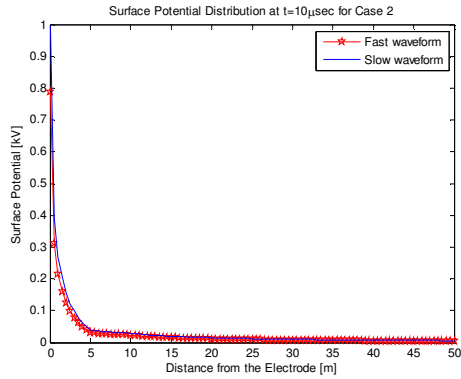


Fig. 20. Surface potential distribution up to a distance of 50m for Case 2 at $t=10\ \mu\text{sec}$ under fast and slow impulse waveforms.

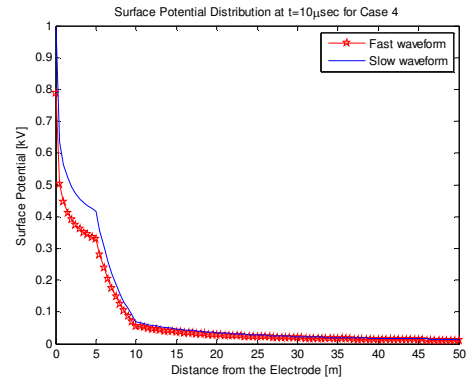


Fig. 24. Surface potential distribution up to a distance of 50m for Case 4 at $t=10\ \mu\text{sec}$ under fast and slow impulse waveforms.

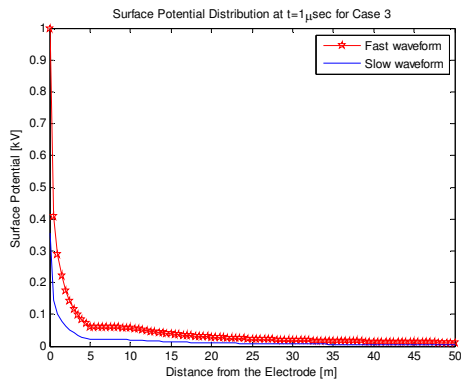


Fig. 21. Surface potential distribution up to a distance of 50m for Case 3 at $t=1\ \mu\text{sec}$ under fast and slow impulse waveforms.

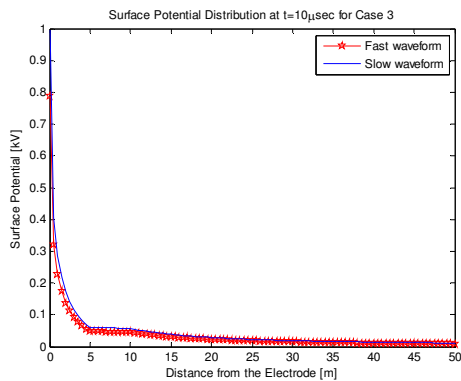


Fig. 22. Surface potential distribution up to a distance of 50m for Case 3 at $t=10\ \mu\text{sec}$ under fast and slow impulse waveforms.

In Figs. 17-24 the potential distribution around the rod is presented for each one of the examined cases for two different waveform impulses (one fast and another slow) for two different time intervals ($t=1\ \mu\text{sec}$ and $t=10\ \mu\text{sec}$). At $1\ \mu\text{sec}$ the fast waveform has reached its maximum value and the slow waveform is ascending. At $10\ \mu\text{sec}$ the slow waveform has reached its maximum value and the fast waveform is decaying from its maximum value. Figs. 17 - 24 show that:

- In those cases where the soil resistivity of the layers closer to the electrode is higher than the resistivity of the next layer (Cases 2 and 3) the potential values for $t=1\ \mu\text{sec}$ and $t=10\ \mu\text{sec}$ are comparable. This shows that the high resistivity region preserves the potential and prevents it from going down to zero (nullifying).
- High values of resistance lead to radical fall of potential regardless of the impulse rise time.

As a result, it is evident that the non-uniformity of the soil and the form of the impulse waveform affect the potential distribution on the surface of the earth. Thus, the multi-layer earth structure must be taken into account when a grounding system is designed.

VI. CONCLUSION

Whenever designing any grounding system the properties of the soil, its heterogeneity as well as the behavior of the grounding system under surge conditions must be taken into consideration. Simulation programs such as OPERA are proved to be an important and useful tool since any earth structure can be simulated. However, there are some

limitations when using OPERA 2d: only simple, axisymmetric grounding systems can be designed. Furthermore, the simulation program does not take into account the non-linear phenomena developed around the grounding system, such as soil ionization. For the simulation of more complicated grounding systems, the 3d version of the program should be used. Moreover, in that case arise some considerations regarding the solvers used by the program. The program solves either the electric or the magnetic problem. That means that either the impedance or the capacitance of the grounding system is taken into consideration.

VII. ACKNOWLEDGMENT

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IX. BIOGRAPHIES

Fani E. Asimakopoulou was born in Athens, Greece, in 1983. She received her Diploma in Electrical and Computer Engineering from the National Technical University of Athens, Greece, in 2006. Currently, she pursues her Ph.D.. Her research interests include high voltages, grounding systems and electromagnetic compatibility. She is a member of the Technical Chamber of Greece.

Vassiliki T. Kontargyri was born in Athens, Greece, in 1978. She received her diploma in electrical and computer engineering and her Ph.D. from National Technical University of Athens in 2002 and 2007, respectively. Her research interests concern high voltage insulators, high voltages, grounding systems and artificial neural networks. She is a member of the Technical Chamber of Greece.

Ioannis F. Gonos was born on May 8, 1970 in Artemisio, Arcadia, Greece. He received his diploma in Electrical Engineering and his Ph.D. from the National Technical University of Athens in 1993 and 2002 respectively. He was a teaching assistant at the Greek Naval Academy and the Technological Education Institute of Athens (1996-2001). He is working at the High Voltage Laboratory of NTUA. His research interests concern grounding systems, insulators, high voltages, electromagnetic immunity and genetic algorithms. He is member of the Technical Chamber of Greece, the IEEE and the CIGRE.

Ioannis A. Stathopoulos was born in Volos, Greece in 1951. He studied in the Faculty of Electrical and Mechanical Engineering of the National Technical University of Athens (1969-1974). He carried out his doctor thesis at the Technical University of Munich (1974-1978). He became teaching assistant at the Technical University of Munich (1974-1978), production engineer in the company "Vianox-Franke" (1979-1980), teaching assistant at the National Technical University of Athens (1979-1983) and thereafter Lecturer (1983-1987), Assistant Professor (1987-1991), Associate Professor (1991-1995) and Professor (since 1995) in the High Voltage Laboratory of the NTUA. He is the author of 8 books and more than 150 papers in scientific journals and conferences proceedings. He is lead assessor of the Hellenic Accreditation Council.