

## EXPERIMENTAL APPLICATION OF THE SIMILARITY THEORY ON SCALED STRIP ELECTRODE CONFIGURATIONS FOR THE CALCULATION OF THE GROUNDING IMPEDANCE

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**Abstract:** Aim of the present work is the rigorous application of the dimensional analysis for the calculation of the impulse impedance of grounding stripes. First, the dimensionless parameters are defined according to the Buckingham Pi-Theorem and the necessary scaling factors for the physical quantities of the problem arise. Stripes embedded in a metal tank filled with saline water are subjected to a lightning impulse produced by an impulse current generator. The waveforms of the injected current and the electrode voltage are recorded for different geometrical scaling factors and conductivities and the impulse impedance is calculated from the oscillogramms. The experimental ratios of the impedance values for the various scaling scenarios are compared to the theoretically expected ratios resulting from the present dimensional analysis. Finally, conclusions are drawn about the accuracy of the proposed method and about factors that influence the experimental results and should be incorporated in the dimensional analysis.

### 1 INTRODUCTION

Experiments on scaled down structures have been adopted in various fields of the engineering science as a way to achieve effective analysis and design of engineering systems. The Similarity Theory is a method applicable on scaled-model experiments, which allows the derivation of scale laws for the physical quantities of a problem, so that the actual values on the full size model can be predicted from measurements on the scaled model. Basis of the theory is the principle of physical similarity between model and prototype, which entails the condition that the equations for original and model differ only by a constant and demands geometrical similarity as a prerequisite [1], [2].

Dimensional analysis is a valuable tool to acquire physical similarity without necessarily having to specify and solve the governing equations of the physical process. By this method we must first identify the physical quantities involved in the problem and then seek dimensionless products formed from them. According to the Buckingham Pi-Theorem [1], [2] all these dimensionless Pi parameters must be identical between physical and scaled model.

In high voltage engineering, the Similarity Theory has been introduced in various areas, such as

lightning protection and grounding system analysis [3], [4], [5]. The first attempt came from Korsuncev for the calculation of surge characteristics of earth electrodes [6]. He described the impulse impedance in terms of two dimensionless arguments and published a collection of relevant data points with an experimentally determined curve. Oettle [7] and Chisholm and Janischewskyj [8] further extended Korsuncev's model by proposing new definitions for the characteristic dimension of the electrode.

### 2 APPLICATION OF THE DIMENSIONAL ANALYSIS

In order to examine the validity of the dimensional analysis, measurements of the grounding impedance of grounding strip electrodes are conducted. For the purposes of the present analysis the following formula proposed by Grcev [9] has been adopted for the calculation of the grounding impedance:

$$A_i = Z/R \quad (1)$$

where:  $Z$  = impulse impedance ( $\Omega$ )  
 $R$  = steady state resistance ( $\Omega$ )  
 $A_i$  = impulse coefficient

This formula takes into consideration the two major physical processes that dominate the

dynamic response of a grounding electrode subjected to a lightning current pulse: 1) the soil ionization, observed as a reduction of the grounding resistance due to the high value of the electric field in the vicinity of the electrode depending on the peak value of the injected current. 2) the limited speed of the lightning pulse propagation along the electrode related to the current pulse front time, thus resulting to an effective discharge of the lightning current to the earth only by a small part of the electrode and a deterioration of the grounding performance. The impulse coefficient taking into consideration both the effective length and ionization phenomena is:

$$A_i = \frac{1}{\sqrt{1+I_m/I_g}} + A - 1, \quad I_g = \frac{E_0 \rho}{2\pi R^2} \quad (2)$$

where:  $E_0$  = the critical ionization field ( $V/m$ )  
 $I_m$  = the peak value of the current (A),  
 $I_g$  = the critical ionization current (A),  
 $A$  = the impulse coefficient without the influence of the ionization, defined as:

$$\begin{aligned} A &= 1 \quad \text{for } (l \leq l_{eff}) \\ A &= \alpha l + \beta \quad \text{for } (l > l_{eff}) \quad (3) \\ l_{eff} &= \frac{1-\beta}{\alpha} \quad \text{is the effective length (m)} \quad (4) \end{aligned}$$

The coefficients  $\alpha$  and  $\beta$  are functions of the resistivity  $\rho$  and the current waveform rise time  $T_1$  (sec). The steady state resistance  $R$  ( $\Omega$ ) of a grounding strip embedded in soil at a depth  $h = 0,5 \div 1,0$  m is given by the formula:

$$R = \frac{\rho}{\pi l} \ln \left( \frac{2l}{d} \right) \quad (5)$$

Where:  $\rho$  = the soil resistivity ( $\Omega \cdot m$ ),  
 $l$  = the length of the grounding strip (m)  
 $d$  = the width of the strip (m)

According to the abovementioned analysis the impulse impedance is a function of the following parameters:

$$Z = f(d, l, \rho, I_m, E_0, T_1) \quad (6)$$

The physical quantities  $d, l, \rho, I_m, E_0, T_1$  are the independent variables of the problem, whereas the impulse impedance  $Z$  is the dependent variable (a total of  $r = 7$  variables that form the problem). Their dimensions on the basis of the unit system  $[M L T I]$ , which consists of the fundamental quantities  $M$  mass,  $L$  length,  $T$  time and  $I$  current are:

$$\begin{aligned} [Z] &= [R] = \Omega = V/A = kg \cdot m^2 \cdot s^{-3} \cdot A^{-2} \Rightarrow \\ [Z] &= [M^1 L^2 T^{-3} A^{-2}] = (1, 2, -3, -2) \quad (7) \\ [d] &= [l] = m \Rightarrow \end{aligned}$$

$$[d] = [l] = [M^0 L^1 T^0 A^0] = (0, 1, 0, 0) \quad (8)$$

$$\begin{aligned} [\rho] &= \Omega \cdot m = kg \cdot m^3 \cdot s^{-3} \cdot A^{-2} \Rightarrow \\ [\rho] &= [M^1 L^3 T^{-3} A^{-2}] = (1, 3, -3, -2) \quad (9) \end{aligned}$$

$$[I_m] = A \Rightarrow [I_m] = [M^0 L^0 T^0 A^1] = (0, 0, 0, 1) \quad (10)$$

$$\begin{aligned} [E_0] &= V/m = kg \cdot m \cdot s^{-3} \cdot A^{-1} \Rightarrow \\ [E_0] &= [M^1 L^1 T^{-3} A^{-1}] = (1, 1, -3, -1) \quad (11) \end{aligned}$$

$$[T_1] = s \Rightarrow [T_1] = [M^0 L^0 T^1 A^0] = (0, 0, 1, 0) \quad (12)$$

Thus the matrix of the dimensions is

$$\begin{bmatrix} 1 & 2 & -3 & -2 \\ 0 & 1 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 1 & 3 & -3 & -2 \\ 0 & 0 & 0 & 1 \\ 1 & 1 & -3 & -1 \\ 0 & 0 & 1 & 0 \end{bmatrix} \quad (13)$$

The number of  $\Pi$ -Parameters to be determined is  $k = r - n = 7 - 4 = 3$ , where  $n = 4$  are the dimensionally independent variables, equal to the rank of the above dimension matrix. Choosing as complete, dimensionally independent subset of variables  $[\rho, I_m, l, T_1]$  the  $\Pi$ -Parameters for the remaining variables arise [1], [2]:

$$\Pi_1 = Z \cdot \rho^{-1} \cdot l^1 = \frac{Z \cdot l}{\rho} \quad (14)$$

$$\Pi_2 = d \cdot l^{-1} = \frac{d}{l} \quad (15)$$

$$\Pi_3 = E_0 \cdot \rho^{-1} \cdot I_m^{-1} \cdot l^2 = \frac{E_0 \cdot l^2}{\rho \cdot I_m} \quad (16)$$

According to the principles of physical similarity, which require that these parameters are equal between model system (M) and full scale system (F), the scale factors for the variables of the problem are defined:

$$\begin{aligned} \Pi_{1F} = \Pi_{1M} &\Rightarrow \frac{Z_F \cdot l_F}{\rho_F} = \frac{Z_M \cdot l_M}{\rho_M} \Rightarrow \frac{Z_F}{Z_M} = \frac{\rho_F}{\rho_M} \cdot \frac{l_M}{l_F} \\ \Rightarrow K_Z = K_\rho \cdot K_l^{-1} &\Rightarrow K_Z = 1/(K_\sigma \cdot K_l) \quad (17) \end{aligned}$$

$$\begin{aligned} \Pi_{2F} = \Pi_{2M} &\Rightarrow \frac{d_F}{l_F} = \frac{d_M}{l_M} \Rightarrow \frac{d_F}{d} = \frac{l_F}{l_M} \\ &\Rightarrow K_d = K_l \quad (18) \end{aligned}$$

$$\begin{aligned} \Pi_{3F} = \Pi_{3M} &\Rightarrow \frac{E_{0F} \cdot l_F^2}{\rho_F \cdot I_{mF}} = \frac{E_{0M} \cdot l_M^2}{\rho_M \cdot I_{mM}} \\ \Rightarrow \frac{E_{0F}}{E_{0M}} \cdot \left( \frac{l_F}{l_M} \right)^2 &= \frac{\rho_F}{\rho_M} \cdot \frac{I_{mF}}{I_{mM}} \Rightarrow K_{E_0} K_l^2 = K_\rho K_{I_m} \\ &\Rightarrow K_{I_m} = K_{E_0} K_l^2 K_\sigma \quad (19) \end{aligned}$$

### 3 EXPERIMENTAL PROCEDURE

The test arrangement is shown in Figure 1. The lightning impulse is produced by a high impulse current generator with maximum stored energy 1.5 kW which generates 8/20 $\mu$ s current waveforms with a peak value up to 25 kA. The output of the generator is connected to electrode

models placed in a rectangular electrolytic tank with dimensions 2m x 1m x 0.5m. The metal walls of the tank are grounded providing a return path for the discharge current.

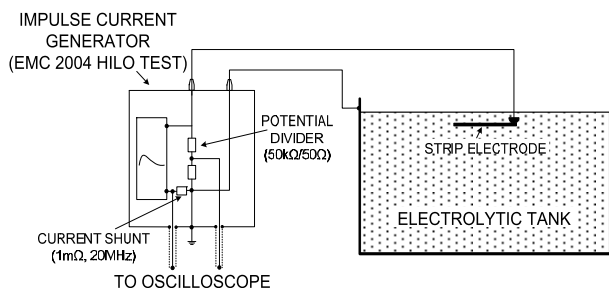


Figure 1: Experimental setup

The strip electrodes are hung with very thin ropes from insulating plastic rods placed across the tank. The hanging is properly adjusted during the measurements in order to achieve a horizontal configuration with the minimum deformation and bend. Three copper stripes, each with dimensions 20cm x 2cm, 30cm x 3cm and 40cm x 4cm are tested. Each electrode under test is located at the centre of the tank at a depth of 2 cm, 3cm and 4cm respectively below the surface of the electrolyte. These depths are chosen to be in accordance with the length and the width of the stripes to maintain geometrical similarity.

The shape and the size of the tank are of great importance because they can distort the field in the vicinity of the model and alter the measurements. Investigations on this matter carried out by previous researchers have shown that a hemispherical tank should have a diameter at least three times the width -or approximately twice the diagonal- of a square grid model [5].

According to this concept, our 2m-long tank can accommodate a strip with up to 60cm length. The rectangular shape of the tank may not exactly simulate a semi-infinite hemispherical earth but it is ideal for the specific strip models which are also rectangular.

The impulse voltage is measured by means of a resistive voltage divider (50kΩ/50Ω) and the injected impulse current by an impulse current shunt (1.0 mΩ) both built in the impulse generator cover. The voltage and current waveforms are recorded with a digital oscilloscope [10].

The salinity of the water is varied in order to achieve the desired conductivity. Taking into consideration that a typical conductivity value for seawater is 5S/m, the tests are conducted for two conductivity levels, 1S/m and 2S/m (the corresponding resistivity values are 1Ωm and 0,5Ωm). The desired injected current is acquired by properly adjusting the generator charging voltage. At each current level five impulses are carried out -to assure the repeatability of the measurements- the voltage and current waveforms are recorded and from the mean value of the electrode peak voltage  $V_m$  and peak current  $I_m$  the impulse impedance is calculated:

$$Z = \frac{V_m}{I_m} \quad (20)$$

The levels of the injected current are defined according to equation (19) so as to allow specific comparisons for each combination of conductivity and electrode dimensions. Thus, the following measurement scenarios are implemented:

- 1) For the same conductivity ( $K_\sigma = K_{E_0} = 1$ ) comparison of the measurements recorded under injected current conditions between electrodes:  $K_{I_m} = K_I^2$ . The scaling law to be verified in this case is:  $K_Z = 1/K_I$
- 2) For the same electrode dimensions ( $K_I = 1$ ) comparison between conductivity cases of the measurements recorded under injected current conditions:  $K_{I_m} = K_{E_0} K_\sigma$ . If we adopt the assumption that the critical ionization field has a constant value ( $K_{E_0} = 1$ ) the above condition is simplified:  $K_{I_m} = K_\sigma$ . The scaling law to be verified in this case is:  $K_Z = 1/K_\sigma$

#### 4 MEASUREMENT RESULTS

In the following tables the recorded current and impedance values are presented along with the deviations  $K_{I_m}$  and  $K_Z$  which are calculated as a percentage of the corresponding theoretically expected scale factors. Especially in the case of the impedance, the theoretical  $K_Z$  is calculated in two ways, either directly as  $K_{Z(a)}$  considering  $K_\sigma = 2$  or from the exact values of the measured conductivities as  $K_{Z(b)}$ . Thus the deviation of the experimental ratio is defined with respect to both theoretical scaling factors i.e. deviation (a) and deviation (b).

**Table 1:** Comparison between the 40x4 and the 30x3 electrodes for  $\sigma=1\text{S/m}$ 

	Theoretical			experimental			$K_{I_m}$ deviation %	theoretical	experimental			$K_Z$ deviation %
	40x4 $I_m$ (A)	30x3 $I_m$ (A)	$K_{I_m}$	40x4 $I_m$ (A)	30x3 $I_m$ (A)	$K_{I_m}$		$K_Z$	40x4 $Z$ ( $\Omega$ )	30x3 $Z$ ( $\Omega$ )	$K_Z$	
	40x4 - 30x3											
	177,8	100	1,778	171,2	102,8	1,665	-6,323	0,750	2,9813	4,0000	0,745	-0,623
	355,6	200		352	198,8	1,771	-0,402		3,0114	4,0563	0,742	-1,015
	533,3	300		535,2	300,8	1,779	0,083		3,0007	4,0293	0,745	-0,701
	711,1	400		710,4	399,2	1,780	0,100		2,9842	3,9980	0,746	-0,476
	888,9	500		880	500	1,760	-1,000		2,9864	4,0000	0,747	-0,455
	1067	600		1062	600	1,770	-0,437		3,0019	3,9667	0,757	0,904
	1244	700		1240	699,2	1,773	-0,243		2,9806	4,0160	0,742	-1,041
	1422	800		1420	800	1,775	-0,156		3,0056	4,0150	0,749	-0,187
	1600	900		1600	897,6	1,783	0,267		2,965	4,0018	0,741	-1,211
	1778	1000		1784	1010	1,766	-0,644		2,9955	3,9762	0,753	0,447
	$K_I = \frac{4}{3} = 1,333$			average values		1,762	-0,876		average values		0,747	-0,436

**Table 2:** Comparison between the 40x4 and the 20x2 electrodes for  $\sigma=1\text{S/m}$ 

	Theoretical			experimental			$K_{I_m}$ deviation %	theoretical	Experimental			$K_Z$ deviation %
	40x4 $I_m$ (A)	20x2 $I_m$ (A)	$K_{I_m}$	40x4 $I_m$ (A)	20x2 $I_m$ (A)	$K_{I_m}$		$K_Z$	40x4 $Z$ ( $\Omega$ )	20x2 $Z$ ( $\Omega$ )	$K_Z$	
	40x4 - 20x2											
	200	50	4	198,4	50,3	3,944	-1,392	0,500	2,9839	5,9141	0,505	0,906
	400	100		400	98,9	4,044	1,112		3,0000	5,8900	0,509	1,868
	600	150		600,8	148,2	4,054	1,350		2,9960	5,8731	0,510	2,024
	800	200		798,4	198	4,032	0,808		3,0110	5,8990	0,510	2,086
	1000	250		1000,4	248,8	4,021	0,523		2,9988	5,8682	0,511	2,206
	1200	300		1200	300	4,000	0,000		2,9867	5,9333	0,503	0,674
	1400	350		1400	349,6	4,005	0,114		3,0000	5,9382	0,505	1,040
	1600	400		1600	400,8	3,992	-0,200		2,9650	5,9581	0,498	-0,471
	2000	500		2004	501,6	3,995	-0,120		2,9940	5,9729	0,501	0,253
	2400	600		2404	601,6	3,996	-0,100		2,9784	5,9840	0,498	-0,456
	2800	700		2800	699,2	4,005	0,114		2,9771	5,9382	0,501	0,271
	$K_I = \frac{4}{2} = 2$			average values		1,762	-0,876			average values		0,505

**Table 3:** Comparison between the 30x3 and the 20x2 electrodes for  $\sigma=1\text{S/m}$ 

	Theoretical			experimental			$K_{I_m}$ deviation %	theoretical	experimental			$K_Z$ deviation %
	30x3 $I_m$ (A)	20x2 $I_m$ (A)	$K_{I_m}$	30x3 $I_m$ (A)	20x2 $I_m$ (A)	$K_{I_m}$		$K_Z$	30x3 $Z$ ( $\Omega$ )	20x2 $Z$ ( $\Omega$ )	$K_Z$	
	30x3 - 20x2											
	225	100	2,25	224,8	98,9	2,273	1,022	0,667	4,0641	5,8900	0,690	3,499
	450	200		448	198	2,263	0,561		4,0000	5,8990	0,678	1,712
	675	300		675,2	300	2,251	0,030		3,9751	5,9333	0,670	0,495
	900	400		897,6	400,8	2,240	-0,466		4,0018	5,9581	0,672	0,748
	1125	500		1130	501,6	2,253	0,124		4,0000	5,9729	0,670	0,454
	1350	600		1350	601,6	2,244	-0,266		4,0000	5,9840	0,668	0,267
	1575	700		1576	699,2	2,254	0,178		3,9645	5,9382	0,668	0,143
	$K_I = \frac{3}{2} = 1,5$			average values		2,254	0,169			average values		0,674

**Table 4:** Comparison between conductivities  $\sigma=2\text{S/m}$  and  $\sigma=1\text{S/m}$  for the 20x2 electrode

20x2	theoretical			experimental			$K_{I_m}$ deviation %	theoretical	experimental			$K_Z$ deviation ( $\alpha$ ) (%)	$K_Z$ deviation ( $\beta$ ) (%)
	$\sigma=2$ $I_m$ (A)	$\sigma=1$ $I_m$ (A)	$K_{I_m}$	$\sigma=2$ $I_m$ (A)	$\sigma=1$ $I_m$ (A)	$K_{I_m}$			$\sigma=2$ $Z$ ( $\Omega$ )	$\sigma=1$ $Z$ ( $\Omega$ )	$K_Z$		
	100	50	2	99,4	50,32	1,975	-1,232	$K_{Z(a)}$  <b>0,500</b>	2,8893	5,9141	0,489	-2,291	-1,834
200	100	198,8		98,88	2,011	0,526	2,9537		5,8900	0,501	0,297	0,765	
300	150	302		148,2	2,038	1,889	2,9563		5,8731	0,503	0,672	1,142	
400	200	401,6		198	2,028	1,414	2,9681		5,8990	0,503	0,632	1,102	
500	250	500		248,8	2,010	0,482	2,9280		5,8682	0,499	-0,207	0,259	
600	300	599,2		300	1,997	-0,133	2,9306		5,9333	0,494	-1,217	-0,755	
700	350	700		349,6	2,002	0,114	2,9200		5,9382	0,492	-1,654	-1,194	
800	400	800		400,8	1,996	-0,200	2,9300	5,9581	0,492	-1,646	-1,187		
900	450	899,2		451,2	1,993	-0,355	2,9493	5,9929	0,492	-1,574	-1,114		
1000	500	1000		501,6	1,994	-0,319	2,9480	5,9729	0,494	-1,287	-0,826		
1200	600	1200		601,6	1,995	-0,266	2,9400	5,9840	0,491	-1,739	-1,280		
1400	700	1400		699,2	2,002	0,114	2,9257	5,9382	0,493	-1,461	-1,001		
measured $K_\sigma = 2,008$				average values		<b>2,003</b>	<b>0,170</b>	average values			<b>0,497</b>	<b>-0,956</b>	<b>-0,494</b>

**Table 5:** Comparison between conductivities  $\sigma=2\text{S/m}$  and  $\sigma=1\text{S/m}$  for the 30x3 electrode

30x3	theoretical			experimental			$K_{I_m}$ deviation %	theoretical	experimental			$K_Z$ deviation ( $\alpha$ ) (%)	$K_Z$ deviation ( $\beta$ ) (%)
	$\sigma=2$ $I_m$ (A)	$\sigma=1$ $I_m$ (A)	$K_{I_m}$	$\sigma=2$ $I_m$ (A)	$\sigma=1$ $I_m$ (A)	$K_{I_m}$			$\sigma=2$ $Z$ ( $\Omega$ )	$\sigma=1$ $Z$ ( $\Omega$ )	$K_Z$		
	200	100	2	200,8	102,8	1,953	-2,335	$K_{Z(a)}$  <b>0,500</b>	2,0438	4,0000	0,511	2,191	3,146
400	200	400,8		198,8	2,016	0,805	2,0200		4,0563	0,498	-0,405	0,526	
450	225	450,4		224,8	2,004	0,178	2,0142		4,0641	0,496	-0,877	0,049	
600	300	601,6		300,8	2,000	0,000	2,0246		4,0293	0,502	0,495	1,434	
800	400	800		399,2	2,004	0,200	2,0200		3,9980	0,505	1,051	1,995	
900	450	899,2		448	2,007	0,357	1,9973		4,0000	0,499	-0,133	0,800	
1000	500	1000		500	2,000	0,000	2,0080	4,0000	0,502	0,400	1,338		
1200	600	1200		600	2,000	0,000	2,0033	3,9667	0,505	1,008	1,952		
1350	675	1350		675,2	1,999	-0,030	2,0089	3,9751	0,505	1,073	2,018		
1400	700	1400		699,2	2,002	0,114	2,0029	4,0160	0,499	-0,257	0,676		
1800	900	1800		897,6	2,005	0,267	2,0156	4,0018	0,504	0,733	1,674		
measured $K_\sigma = 2,020$				average values		<b>2,000</b>	<b>0,007</b>	average values			<b>0,502</b>	<b>0,480</b>	<b>1,419</b>

## 5 CONCLUSIONS

A comparison of the results shows an overall satisfying performance of the theory. The first measurement scenario (same conductivity) has yielded very small impedance deviations, not exceeding 1%. The deviations at the second measurement scenario (same electrode) are slightly higher. For all the electrodes the deviations  $K_{Z(b)}$  are smaller than the deviations  $K_{Z(a)}$ . This proves that the exact value of the conductivity must be taken into account.

These deviations can be justified by the limited accuracy of the instruments (generator and oscillograph). A differential voltage probe and a current monitor would increase the accuracy of the measurements.

The effect of the  $K_{I_m}$  deviation (not more than 1%) was proven to be negligible. Special attention must be given at the exact adjustment of the conductivity so that the comparison of the results can be done on the proper basis.

**Table 6:** Comparison between conductivities  $\sigma=2\text{S/m}$  and  $\sigma=1\text{S/m}$  for the 40x4 electrode

40x4	theoretical			experimental			$K_{I_m}$ deviation %	theoretical	experimental			$K_z$ deviation (a) %	$K_z$ deviation (b) %
	$\sigma=2$ $I_m$ (A)	$\sigma=1$ $I_m$ (A)	$K_{I_m}$	$\sigma=2$ $I_m$ (A)	$\sigma=1$ $I_m$ (A)	$K_{I_m}$		$K_{Z(a)}$	$\sigma=2$ $Z$ ( $\Omega$ )	$\sigma=1$ $Z$ ( $\Omega$ )	$K_z$		
	200	100		202,4	98,2	2,061	3,055		1,5771	2,9369		0,537	7,399
	400	200	2	400,8	198,4	2,020	1,008	<b>0,500</b>	1,5888	2,9839	0,532	6,494	2,215
	600	300		600,8	298	2,016	0,805		1,5872	3,0121	0,527	5,390	1,156
	800	400		800	400	2,000	0,000		1,6000	3,0000	0,533	6,667	2,381
	1000	500		998,4	499,2	2,000	0,000		1,5960	3,0008	0,532	6,372	2,098
	1200	600		1200	600,8	1,997	-0,133	<b>0,521</b>	1,5867	2,9960	0,530	5,919	1,663
	1400	700		1400	700,8	1,998	-0,114		1,5714	2,9795	0,527	5,484	1,246
	1600	800		1600	798,4	2,004	0,200		1,5875	3,0110	0,527	5,446	1,209
	2000	1000		2000	1000,4	1,999	-0,040		1,5900	2,9988	0,530	6,042	1,782
	measured $K_\sigma = 1,919$			average values		<b>2,011</b>	<b>0,531</b>		average values		<b>0,531</b>	<b>6,135</b>	<b>1,870</b>

An initial proposal for future improvement of this work would be the processing of the measurements with an alternative definition of the impulse impedance that takes ionization consideration by accounting for the instantaneous voltage value at the current peak.

In the present analysis the critical ionization field was considered to have a steady value. The influence of this parameter and its possible dependence from the conductivity highlighted by other researchers should be investigated through experiments at different salinity levels. Higher conductivity levels would make the effect of the ionization procedure more obvious. In the presented measurements the recordings stopped at specific charging voltage levels where sparks on the tank base were observed. Thus the injected current -which determines the ionization procedure- was limited by the depth of the tank. Finally, other electrode configurations would enable the investigation of the alternative definitions given for the characteristic dimension of the electrode suggested by Korsuncev, Popolansky and Oettle.

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